

THE INTERGALACTIC NEWTONIAN GRAVITATIONAL FIELD AND THE SHELL THEOREM

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SUMMARY: The release of the 2MASS Redshift Survey (2MRS) with its 44599 galaxies allows the deduction of their masses in nearly complete sample. A cubic box with side of 37 Mpc containing 2429 galaxies is extracted and the Newtonian gravitational field is evaluated both at the center of the box as well as in 101 x 101 x 101 grid points of the box. The obtained results are then discussed in the light of the shell theorem which states that inside of a sphere the gravitational field is zero.

Key words. methods: statistical – cosmology: observations – large-scale structure of Universe

1. INTRODUCTION

The determination of the gravitational field in cosmology varies between the Newton law and various types of modifications to this law. The reference formula is the Newtonian force:

$$F = -G \frac{mM}{r^2}, \quad (1)$$

where G is the gravitational constant, M the first mass, m the second mass and r the distance between the two masses. The enormous progress in the observations of the spatial distribution of galaxies points toward a cellular structure, i.e. the galaxies are situated on the surfaces of bubbles, rather than to be aggregated in a random structures (see Coil 2012). In the limiting case in which all the galaxies were situated on the surfaces of spheres, the gravitational forces should be zero due to the shell theorem or nearly zero due to the fact that the galaxies are distributed in a discrete way rather than in a continuous way. This paper describes in Section 2 two astronomical catalogs which allow to calibrate the size of the cosmic voids. Section 3 is devoted to the study

of the photometric properties of a nearly spherical distribution of galaxies and to a careful analysis of completeness connected with the selected astronomical catalog. Section 4 contains the evaluation of the Newtonian gravitational field in a box of 37 Mpc in which the boundary conditions are properly evaluated. Section 5 reports on a comparison between three ideal structures and a real void as extracted from a slices oriented catalog.

2. OBSERVATIONS

This section processes the Sloan Digital Sky Survey Data Release 7 (SDSS DR7), see Abazajian et al. (2009) and the 2MASS Redshift Survey (2MRS), see Huchra et al. (2012).

2.1. Observed statistics of the voids

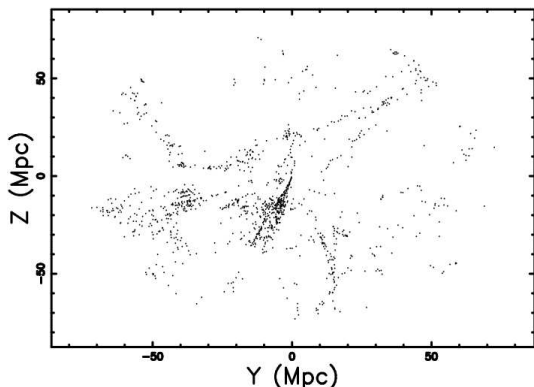
The distribution of the effective radii between the galaxies of SDSS DR7 has been reported in Pan et al. (2012). This catalog contains 1054 voids and Table 1 reports their basic statistical parameters.

Table 1. The statistical parameters of the effective radii in SDSS DR7.

parameter	value
mean	$18.23h^{-1}$ Mpc
variance	$23.32h^{-2}$ Mpc ²
standard deviation	$4.82h^{-1}$ Mpc
kurtosis	0.038
skewness	0.51
maximum value	$34.12h^{-1}$ Mpc
minimum value	$9.9h^{-1}$ Mpc

2.2. The 2MASS

The 2MASS is a catalog of galaxies which contains observations in the near-infrared J, H and K-bands (1-2.2 μm) and therefore detects the galaxies in the so called "Zone of Avoidance", see Jarrett (2004) and Crook et al. (2007). At the moment of writing, the 2MRS consists of 44599 galaxies with redshift in the interval $0 \leq z \leq 0.09$ (Huchra et al. 2012). The catalog gives the galactic latitude, the galactic longitude and the expansion velocity; from these three parameters it is possible to deduce the Cartesian coordinates, X , Y and Z expressed in Mpc. Fig. 1 reports on a cut of a given thickness of 2MRS where Δ expresses the thickness of the cut and N_G the number of selected galaxies.


Fig. 1. Cut of the 3D spatial distribution of 2MRS in the $X = 0$ plane when $\Delta = 10$ Mpc, the squared box has a side of 148 Mpc and $N_G = 1244$.

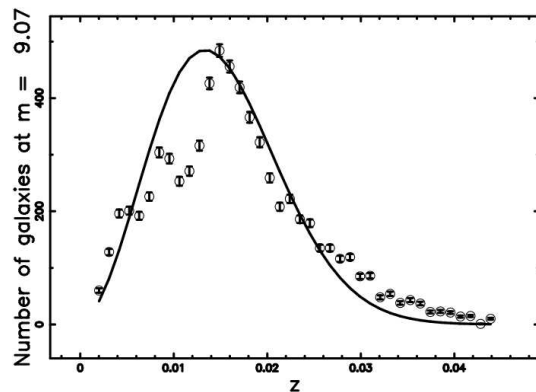
3. PHOTOMETRIC PROPERTIES

This section reviews the photometric maximum in the framework of the luminosity function for galaxies and the Malmquist bias which fixes the concept of a complete sample. A model for the luminosity of galaxies is the Schechter function $\Phi(L; L^*, \alpha, \Phi)$ where α denotes the slope for low values of L , L^* is the characteristic luminosity, and Φ^* is a normalization, see Eq. (55) in Zaninetti (2010b). This function was suggested by Schechter (1976) and the distribution in absolute magnitude $\Phi(M; M^*, \alpha, \Phi)$ can be

found in Eq. (56) in Zaninetti (2010b) where M^* is the characteristic magnitude as derived from the data. The parameters of the Schechter function concerning the 2MRS as well as the bolometric luminosity $M_{K_S}^\odot$ can be found in Cole et al. (2001) and are reported in Table 2.

Table 2. The parameters of the Schechter function and bolometric magnitude for the 2MRS in the $K_s - \text{band}$.

parameter	2MRS
$M^* - 5 \log_{10} h$ [mags]	(-23.44 ± 0.03)
α	-0.96 ± 0.05
$\Phi^* [h^3 \text{ Mpc}^{-3}]$	$((1.08 \pm 0.16)10^{-2})$
$M_{K_S}^\odot$	3.39
h	0.7


Fig. 2. The galaxies of the 2MRS with $8.48 \leq m \leq 10.44$ or $1202409 \frac{L_{* \odot}}{\text{Mpc}^2} \leq f \leq 7267112 \frac{L_{* \odot}}{\text{Mpc}^2}$ organized in frequencies versus heliocentric redshift, (empty circles); the error bar is given by the square root of the frequency. The maximum frequency of observed galaxies is at $z = 0.015$. The full line is the theoretical curve generated by $\frac{dN}{d\Omega dz df}(z; z_{\text{crit}}, c, H_0)$. In this plot, $M_{K_S}^\odot = 3.39$, $h = 0.7$, $M^* = -24.87$, $\alpha = -0.98$, $\Phi^* = 0.0037$, $\chi^2 = 721$ and the number of bins 40.

The number of galaxies at a given flux f as a function of the redshift z , $\frac{dN}{d\Omega dz df}(z; z_{\text{crit}}, c, H_0)$, are given by Eq. (1.104) of Padmanabhan (1996) or by Eq. (6) of Zaninetti (2010b), where $d\Omega$, dz , and df are the differentials of the solid angle, the red-shift, and the flux, respectively, z_{crit} is a parameter, H_0 the Hubble constant and c is the velocity of light. The number of galaxies at a given flux has a maximum at $z = z_{\text{max}}(z_{\text{crit}}, \alpha)$, see Eq. (8) in Zaninetti (2010b). Fig. 2 reports on the number of observed galaxies in the 2MRS catalog at a given apparent magnitude and the theoretical curve as represented by $\frac{dN}{d\Omega dz df}(z; z_{\text{crit}}, c, H_0)$. The merit function χ^2 can be computed as:

$$\chi^2 = \sum_{j=1}^n \left(\frac{n_{\text{theo}}(z) - n_{\text{astr}}(z)}{\sigma_{n_{\text{astr}}(z)}} \right)^2, \quad (2)$$

where n is the number of data, the two indices theo and astr stand for theoretical and astronomical, respectively and $\sigma_{n_{\text{astr}}(z)}^2$ is the variance of the astronomical number of data; the obtained value is reported in the caption of Fig. 2.

The total number of galaxies in the 2MRS as function of z is reported in Fig. 3 as well as the theoretical curve represented by the numerical integration of $\frac{dN}{d\Omega dz df}(z; z_{\text{crit}}, c, H_0)$.

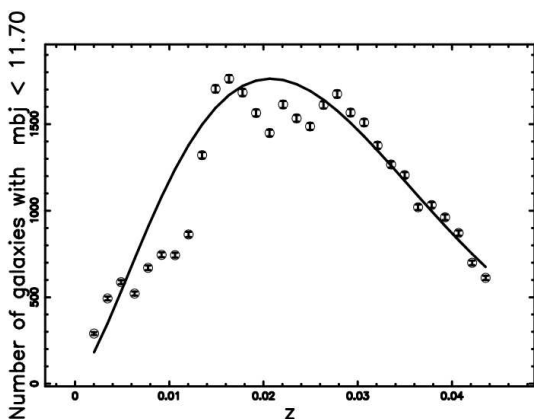


Fig. 3. All the galaxies of the 2MRS with $m_{ks} < 11.75$ organized in frequencies versus heliocentric redshift (empty circles); the error bar is given by the square root of the frequency. The maximum frequency of all observed galaxies is at $z = 0.017$. The full line is the theoretical curve generated by $\frac{dN}{d\Omega dz df}(z; z_{\text{crit}}, c, H_0)$. In this plot $M_{K_S}^{\odot} = 3.39$, $h = 0.7$, $M^* = -23.97$, $\alpha = -0.96$, $\Phi^* = 0.0037$, $\chi^2 = 1267$ and the number of bins is 30.

The mass of a galaxy can be evaluated once the mass luminosity ratio R , is given, by:

$$R = \left\langle \frac{M}{L} \right\rangle. \quad (3)$$

Some values of R are now reported: $R \leq 20$ by Kiang (1961) and Persic and Salucci (1992), $R = 20$ by Padmanabhan (1996) and $R = 5.93$ by van der Marel (1991). The Malmquist bias, (Malmquist 1920, 1922), was originally applied to the stars, and later on to the galaxies by Behr (1951). The observable absolute magnitude $M_L(m_L; z, H_0)$ as a function of the limiting apparent magnitude m_L , is given by Eq. (51) in Zaninetti (2010b). The bias predicts, from a theoretical point of view, an upper limit for the maximum absolute magnitude which can be observed in a catalog of galaxies characterized by a given limiting magnitude and Fig. 4 reports such a curve as well as the galaxies of the 2MRS.

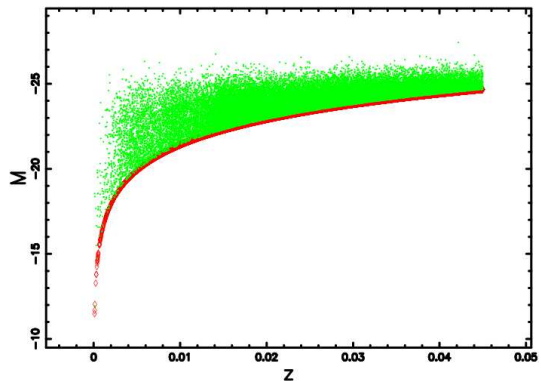


Fig. 4. The absolute magnitude M of 36464 galaxies belonging to the 2MRS when $M_{K_S}^{\odot} = 3.39$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ (green points). The upper limit theoretical curve as represented by $M_L(m_L; z, H_0)$ is reported as the red thick line when $m_L = 11.75$.

The limiting magnitude of the 2MRS is $m_L = 11.75$ and therefore the 2MRS is complete for $z \leq 0.0025$. For values of z greater than this value the observed sample is not complete and we can introduce the efficiency $\epsilon_s(z; M_{\text{max}}, M_{\text{min}}, m_L, c, h)$ where M_{max} and M_{min} are the maximum and minimum absolute magnitudes of the considered catalog and $h = H_0/100$, see Eqs. (51-53) in Zaninetti (2010b). As an example, when $z \approx 0.017$ the sample covers the 52% of the range in absolute magnitude.

4. THE GRAVITATIONAL FIELD

We now explore a connection with the shell theorem. Let us consider a spherically symmetric surface of radius a on which a total mass M is distributed in a uniform way. The mass per unit area μ is:

$$\mu = \frac{M}{4\pi a^2}. \quad (4)$$

The force inside the spherical surface is

$$\Gamma = 0, \quad r < a, \quad (5)$$

see Eq. (11.24) in Alonso and Finn (1992) and, therefore, the shell theorem can be formulated: "A uniform shell of matter exerts no gravitational force on a particle situated inside a shell". We now compute the force in the center ($x=0, y=0, z=0$) of a hemisphere which resides on the positive z -axis. The vectorial intensity of the field is:

$$d\Gamma = G \frac{dm}{a^2}, \quad (6)$$

being

$$dm = \mu a^2 d\Omega, \quad (7)$$

where $d\Omega = \sin\theta d\theta d\phi$ is the solid angle with $0 \leq \phi \leq 2\pi$ and $0 \leq \theta \leq \frac{\pi}{2}$. The three Cartesian components of the field in the 3D case are:

$$\begin{aligned}\Gamma_z &= \int_0^{2\pi} d\phi \int_{\frac{\pi}{2}}^0 G \mu \sin\theta \cos\theta d\theta, \\ \Gamma_x &= \int_0^{2\pi} d\phi \int_{\frac{\pi}{2}}^0 G \mu \sin\theta \sin\theta \cos\phi d\theta, \\ \Gamma_y &= \int_0^{2\pi} d\phi \int_{\frac{\pi}{2}}^0 G \mu \sin\theta \sin\theta \sin\phi d\theta.\end{aligned}\quad (8)$$

The integration gives in the 3D case for the three forces at the center:

$$\begin{aligned}\Gamma_z &= \frac{GM}{2a^2}, \\ \Gamma_x &= \Gamma_y = 0.\end{aligned}\quad (9)$$

We now consider the 2D case of mass concentrated on a half circle of radius a situated on the positive y -axis where, now, the mass per unit length μ_{2D} is:

$$\mu_{2D} = \frac{M}{\pi a}.\quad (10)$$

The 2D vectorial intensity of the field is:

$$d\Gamma = G \frac{dm}{a^2},\quad (11)$$

with:

$$dm = \mu_{2D} a d\theta,\quad (12)$$

where $0 \leq \theta \leq \pi$. The 2D Cartesian gravitational components of the force at the center ($x=0, y=0$) are:

$$\begin{aligned}\Gamma_y &= G \frac{\mu_{2D}}{a} \int_0^\pi \sin\theta d\theta, \\ \Gamma_x &= G \frac{\mu_{2D}}{a} \int_0^\pi \cos\theta d\theta.\end{aligned}\quad (13)$$

The integration of the 2D case gives:

$$\begin{aligned}\Gamma_y &= \frac{2GM}{\pi a^2}, \\ \Gamma_x &= 0.\end{aligned}\quad (14)$$

At the moment of writing the Committee on Data for Science and Technology (CODATA) recommends:

$$G = (6.67384 \pm 0.00080) \times 10^{-11} \frac{m^3}{kg s^2},\quad (15)$$

Table 3. 3D Gravitational forces expressed in $\frac{Mpc M_{gal}}{yr^2}$ at the center of a 3D box of side 37×2 Mpc when $R=6$ and theoretical 3D formula (9). At $z=0.008$, the efficiency of the sample is $\approx 70.6\%$.

Environment	Γ_x	Γ_y	Γ_z	Γ
real structure	$-9.77 \cdot 10^{-6}$	$-1.53 \cdot 10^{-5}$	$-3.04 \cdot 10^{-6}$	$1.84 \cdot 10^{-5}$
half sphere	0	0	$2.55 \cdot 10^{-5}$	$2.55 \cdot 10^{-5}$

see Mohr et al. (2008). Before to continue we express the Newtonian gravitational constant in the following units: length in Mpc, mass in M_{gal} which is $10^{11} M_{\odot}$ and yr^2 which are $10^8 yr$

$$G = 4.4997510^{-6} \frac{Mpc^3}{M_{gal} yr^2}.\quad (16)$$

The two formulae (9) in 3D and (14) in 2D represent a useful reference to test a numerical code and to fix the range of variability of the gravitational field. According to our 3D theory, the gravitational field at the center of the cosmic voids varies between the minimum value of zero (shell theorem) and a maximum value:

$$\Gamma_z = \frac{GM}{2R^2} = 6.769 \cdot 10^{-9} N \frac{Mpc M_{gal}}{yr^2},\quad (17)$$

where N is the number of galaxies in the spherical shell surrounding the cosmic void having mass $M = M_{gal}$, and $\bar{R}=18.23$ Mpc is the average radius of the cosmic voids.

We are now ready to process the 2MRS data and we associate to each galaxy, as reported in Fig. 1, a mass given by Eq. 3. The three components of the gravitational field are reported in Table 3.

From a careful analysis of Table 3 it is possible to conclude that the gravitational field is greater than zero but smaller with respect to the case in which all the galaxies reside on a half sphere of radius equal to the averaged radius of the sample. Fig. 5 reports on a slice at the middle of a smaller box.

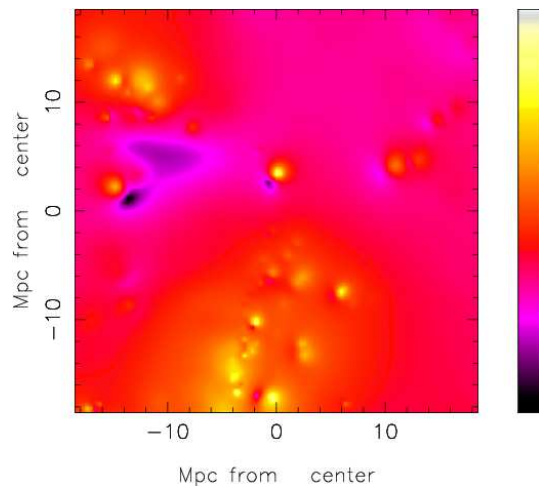


Fig. 5. Cut-map of the 3D gravitational field of 2MRS when $R=6$. In order to have periodic boundary conditions the side of the box is 37 Mpc.

The spatial displacement of the 3D grid $\Gamma(i, j, k)$ which represents the absolute value of the gravitational field can be visualized through the iso-density contours and, as an example, we considered a $101 \times 101 \times 101$ grid. In order to do so, the maximum value $\Gamma_n(i, j, k)_{\max}$ and the minimum value $\Gamma_n(i, j, k)_{\min}$ should be extracted from the three-dimensional grid. A value of this grid can be fixed by the following equation:

$$\Gamma_n(i, j, k)_{\text{chosen}} = \Gamma_n(i, j, k)_{\min} + (\Gamma_n(i, j, k)_{\max} - \Gamma_n(i, j, k)_{\min}) \times \text{coef}, \quad (18)$$

where coef is a parameter comprised between 0 and 1. This iso-surface rendering of the gravitational field is reported in Fig. 6; the Euler angles characterizing the point of view of the observer are also reported.

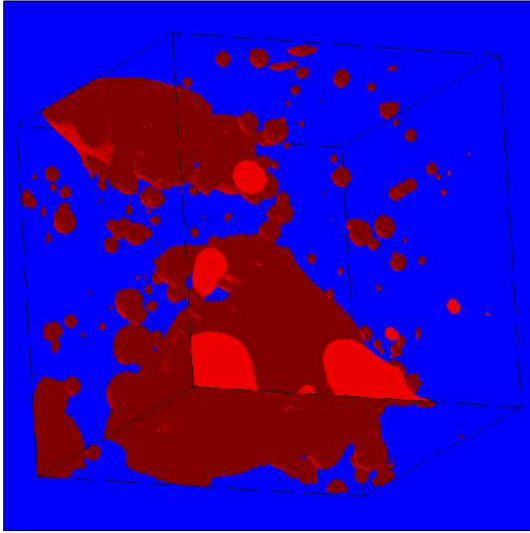


Fig. 6. Iso-surface of the logarithm of the 3D gravitational field of 2MRS when $R=6$ and $\text{coef} = 0.43$. The orientation of the figure is characterized by the Euler angles, which are $\Phi=30^\circ$, $\Theta=30^\circ$ and $\Psi=30^\circ$.

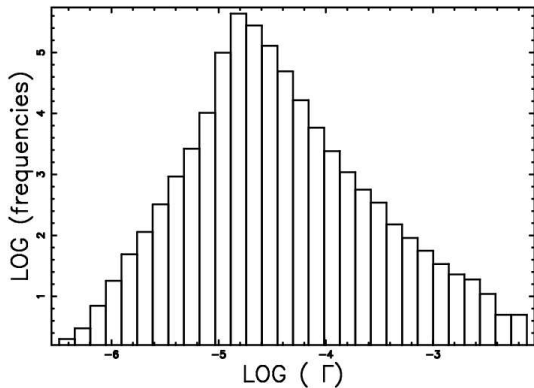


Fig. 7. Decimal logarithmic histogram (step-diagram) of the values of the gravitational field evaluated in $101 \times 101 \times 101$ points.

Another interesting quantity to plot is the statistics of values of already defined spatial grid Γ_n which holds $101 \times 101 \times 101$ values of gravitational field, see Fig. 7.

From this histogram it is possible to conclude that 90% of the space has a gravitational field within $3.24 \cdot 10^{-7} \frac{\text{Mpc} M_{\text{gal}}}{\text{yr}^2} \leq \Gamma_n \leq 3.16 \cdot 10^{-5} \frac{\text{Mpc} M_{\text{gal}}}{\text{yr}^2}$.

5. THE VORONOI SIMULATION

The Poisson Voronoi tessellation (PVT) is a useful tool to explore the spatial clustering of galaxies. The filaments of galaxies visible in the slices-type catalogs are due to the intersection between a plane and the PVT network of faces as the first approximation. An improvement can be obtained by coding the intersection between the slice of a given opening angle and the PVT network of faces (Zaninetti 2006, 2010a). As an example, Fig. 8 reports both the CFA2 slice as well as the simulated slice.

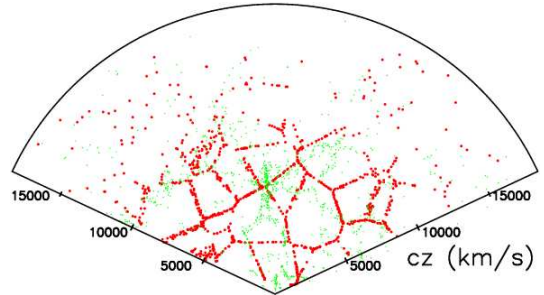


Fig. 8. Polar plot of real galaxies (green points) belonging to the second CFA2 redshift catalog and the simulated galaxies in the PVT framework (red points). More details can be found in Zaninetti (2006).

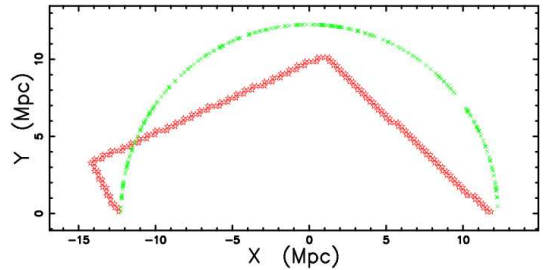


Fig. 9. Half circle with $a=12.25$ Mpc (green stars) and half irregular polygon (red squares). The number of unit masses, $1 M_{\text{gal}}$, is 455.

We now test formula (14) in a discrete environment rather than in the continuous case. The test now calculates the two forces, Γ_x and Γ_y , in the center of the circle and in a 2D irregular Voronoi polygon generated by PVT which has the same averaged radius of the circle and center occurring in

the same location of the generating seed. The half Voronoi polygon and the half circle are displayed in Fig. 9 and the two forces, Γ_x and Γ_y , in Table 4.

Table 4. Gravitational forces expressed in $\text{Mpc}M_{\text{gal}}/\text{yr}^2$ in the comparison between half circle, 2D formula (14), and half irregular polygon. The parameters are $M_{\text{gal}}=455$ and $a = 12.25453$ Mpc.

Environment	Γ_x	Γ_y
Half circle -theory	0	$8.67 \cdot 10^{-6}$
Half circle -numeric	$-4.16 \cdot 10^{-8}$	$8.81 \cdot 10^{-6}$
Half polygon-numeric	$1.37 \cdot 10^{-6}$	$1.36 \cdot 10^{-5}$

We are now ready to process a real void and our attention is focused on a CFA2 slice shown in Fig. 10.

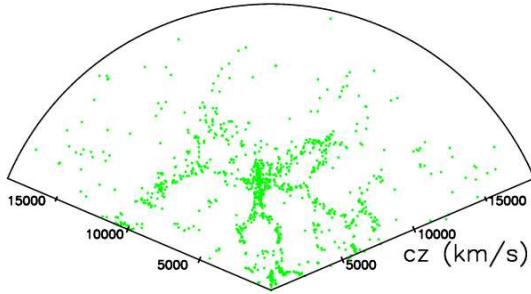


Fig. 10. Polar plot of the real galaxies (green points) belonging to the second CFA2 redshift catalog.

A real void is extracted and the averaged radius of the galaxies on the boundary of that void is computed, see Fig. 11.

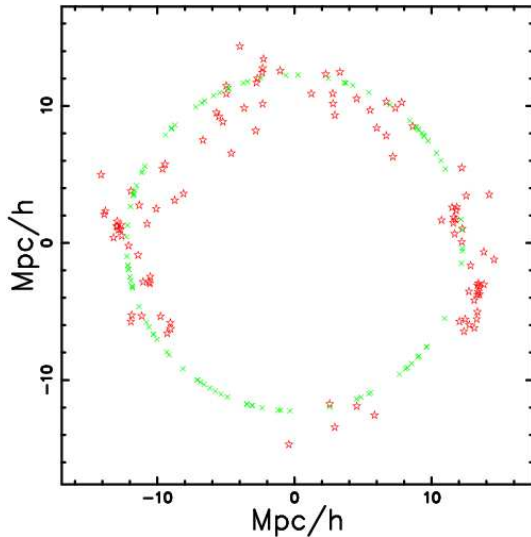


Fig. 11. Circle with $a= 12.25$ Mpc (green stars) and real void extracted from a CFA2 slice (red squares). The number of galaxies with unit mass is 101.

The forces in the x and y direction are then computed and reported in Table 5.

Table 5. Gravitational forces expressed in $\frac{\text{Mpc}M_{\text{gal}}}{\text{yr}^2}$ for the comparison between a circle and a real CFA2 void. The parameters are $M_{\text{gal}}=101$ and $a = 12.25453$ Mpc.

Environment	Γ_x	Γ_y
circle-theory	0	0
half circle-theory	0	$1.92 \cdot 10^{-6}$
real void-numeric	$-1.36 \cdot 10^{-7}$	$8.3 \cdot 10^{-7}$

The presence of both a discrete number of galaxies and a not exactly symmetric displacement of the galaxies produces gravitational forces that take a finite value rather than zero. It is interesting to point out that Γ_y , due to the galaxies on the boundary of the real void, is smaller than the theoretical value as given by the half circle which represents a maximum theoretical value.

6. CONCLUSIONS

The masses of the galaxies can be deduced starting from the luminosities in the framework of the mass luminosity ratio R . The spatial distribution of the masses of the galaxies allows the computation of the Newtonian gravitational forces acting on the unit mass. As a reference for the evaluation of the forces the 2D and 3D shell theorem is analyzed. The evaluation of forces at the center of the box allows to conclude that the forces are smaller with respect to the mass concentrated on a half sphere of radius equal to the averaged radius of the selected sample of galaxies, but bigger than zero due to the fact that the distribution of the galaxies is discrete rather than continuous. A careful analysis of a cubic box having sides of 37 Mpc allows to state that 90 % of the space has gravitational forces around the average value of $2.1 \cdot 10^{-5} \frac{\text{Mpc}M_{\text{gal}}}{\text{yr}^2}$.

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МЕЂУГАЛАКТИЧКО ЊУТНОВО ГРАВИТАЦИОНО ПОЉЕ
И ТЕОРЕМА О СФЕРНО-СИМЕТРИЧНОЈ ЉУСЦИ

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Оригинални научни рад

Објављивање прегледа неба 2MASS Redshift Survey (2MRS) који садржи 44599 галаксија, омогућује изучавање маса галаксија у скоро целокупном узорку. Издвојена је коцка са страницама од 37 Мрс која садржи 2429 галаксија и израчунато је Њутново гравитационо поље, како у центру коцке, тако и на

мрежи од $101 \times 101 \times 101$ тачака унутар ње. Добијени резултати су затим дискутовани у светлу теореме о сферно-симетричној щусци (шупљој лопти), која постулира да је гравитационо поље једнако нули у унутрашњости такве щуске.